ON LIMIT CYCLES APPEARING BY POLYNOMIAL PERTURBATION OF DARBOUXIAN INTEGRABLE SYSTEMS

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ABSTRACT. We prove an existential finiteness result for integrals of rational 1-forms over the level curves of Darbouxian integrals.

1. LIMIT CYCLES BORN BY PERTURBATION OF INTEGRABLE SYSTEMS

1.1. **Poincaré–Pontryagin integral.** Limit cycles (isolated periodic trajectories) of polynomial planar vector fields can be produced by perturbing integrable systems which have nested continuous families of non-isolated periodic trajectories. The number and position of limit cycles born in such perturbations is determined by the number and position of *isolated zeros* of the Poincaré–Pontryagin integral of the dissipation form along the closed periodic orbits of the non-perturbed integrable vector field. The *infinitesimal Hilbert problem* is to place an upper bound for the number of isolated zeros of this integral. This problem was repeatedly formulated as a relaxed form of the Hilbert 16th problem by V. Arnold [1].

Instead of polynomial vector fields, it is more convenient to deal with (singular) foliations of the real plane \mathbb{R}^2 by solutions of rational Pfaffian equations $\theta = 0$, where θ is a 1-form on \mathbb{R}^2 with rational coefficients (since only the distribution of null spaces of the form θ makes geometric sense, one can always replace the rational form by a polynomial one). The foliation \mathcal{F} is (Darbouxian) *integrable*, if the form is *closed*, $d\theta = 0$. An integrable foliation always has a "multivalued" first integral of the form $f(x, y) = \exp r(x, y) \cdot \prod_j p_j(x, y)^{\lambda_j}$, where r(x, y) is a rational function and p_j polynomials in x, y, which are involved in the in general non-rational powers λ_j . A particular case of the integrable foliations consists of Hamiltonian foliations defined by the exact polynomial 1-form $\theta = dh$, $h \in \mathbb{R}[x, y]$.

If $L \subset \mathbb{R}^2$ is a compact smooth leaf (oval) of an integrable foliation $\theta = 0$ belonging to the level curve $\{f = a\} \subset \mathbb{R}^2$, then all nearby leaves of this foliation are also closed and form a continuous family of ovals belonging to the level curves $\{f = t\}, t \in (\mathbb{R}^1, a)$. The ovals from this family, denoted by L_t , are uniquely parameterized by the values of the real variable t varying

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in an open interval $\alpha < t < \beta$. Their union is an annulus $A_{\alpha,\beta}$ bounded by the Hausdorff limits $\overline{L}_{\alpha} = \lim_{t\to\alpha+0} L_t$ and $\overline{L}_{\beta} = \lim_{t\to\beta-0}$. The limits themselves may be not leaves of the foliation, but rather *separatrix polygons* (unions of several leaves and one or more singular points of the foliation \mathcal{F}). The polar locus of the form θ is usually a separatrix polygon.

Consider an one-parameter family of planar real foliations $\mathcal{F}_{\varepsilon}$ defined by the Pfaffian equations

$$\theta + \varepsilon \omega = 0, \quad \text{Poles}(\omega) \subseteq \text{Poles}(\theta), \quad (1.1)$$

with a closed rational 1-form θ and arbitrary rational 1-form ω . The assumption on the poles guarantees that the perturbation (1.1) does not create additional singularities in the annulus $A_{\alpha,\beta}$ (such ω will be called *admissible* for θ). The *Poincaré–Pontryagin integral*¹ associated with this family, is the function

$$t \mapsto I(t) = \oint_{L_t} \omega, \qquad L_t \subseteq \{f = t\} \text{ an oval, } t \in (\alpha, \beta).$$
 (1.2)

The following well-known Poincaré–Pontryagin criterion establishes connections between the integral (1.2) and limit cycles of the system (1.1): If the function I(t) has an isolated zero of multiplicity μ at an interior point $a \in (\alpha, \beta)$, then the perturbed foliation $\mathcal{F}_{\varepsilon}$ for all sufficiently small values of ε and $\delta > 0$ has no more than μ limit cycles in the thin annulus between the ovals $L_{a-\delta}$ and $L_{a+\delta}$. The assertion is reasonably sharp: under certain additional assumptions of nondegeneracy one can produce exactly μ limit cycles.

Thus bounds for the number of isolated zeros for the integral (1.2) translate into bounds for the number of limit cycles for near-integrable systems.² The following conjecture seems to be believed by most experts in the area.

Conjecture 1. For any pair of rational 1-forms (θ, ω) on the real plane \mathbb{R}^2 of degrees n, m respectively, such that $d\theta = 0$ and $\text{Poles}(\omega) \subseteq \text{Poles}(\theta)$, the number of isolated real zeros of the integral (1.2) is bounded by a constant N = N(n, m) depending only on n and m.

Demonstration (or refutation) of this Conjecture constitutes the *infini*tesimal Hilbert Sixteenth problem [1,5]. A separate question concerns computability (and practical computation) of the bound N(n,m).

1.2. Analytic properties of the Poincaré–Pontryagin integral: Hamiltonian versus Darbouxian cases. The function I(t) defined by the integral (1.2) is obviously real analytic inside the real interval (α, β) , and hence the number of isolated zeros is necessarily finite on any compact subinterval

¹The function I(t) is sometimes called Melnikov function or (misleadingly) the Abelian integral, see below.

²However, the above Poincaré–Pontryagin criterion does not apply to limit cycles born from separatrix polygons \overline{L}_{α} , \overline{L}_{β} which bound the annuli of periodic orbits of the integrable systems. This bifurcation requires much deeper analysis that will not be discussed here.

of this interval. Yet this circumstance neither implies any explicit upper bound, nor rules out the accumulation of infinitely many isolated roots to the boundary points α, β of the interval.

Most studies of the function I(t) are based on its analytic continuation from the real interval (α, β) to the complex domain. Already at this step drastic differences between Hamiltonian and general integrable cases occur.

1.2.1. Abelian integrals. If the initial foliation $\mathcal{F} = \{\theta = 0\}$ is Hamiltonian, i.e., if $\theta = dh$, $h \in \mathbb{R}[x, y]$, then ω must also be polynomial and the function I(t) is an Abelian integral, integral of a polynomial 1-form over a cycle on algebraic curve. Such integrals possess the following properties.

(i) The integral I(t) admits analytic continuation on the entire Riemann sphere $\mathbb{C}P^1 = \mathbb{C} \cup \{\infty\}$ as a multivalued function ramified over finitely many points. For a generic Hamiltonian h these points are exactly the (complex) critical values of h.

(ii) I(t) is (a linear combination of coordinates of) a solution of a Fuchsian linear system on $\mathbb{C}P^1$.

(iii) Near each singular point t_0 of the Fuchsian system, I(t) admits a local representation as the finite sum $I(t) = \sum_{j,k} a_{jk}(t) (t - t_0)^{\lambda_j} \ln^k(t - t_0)$ with the coefficients a_{jk} holomorphic at the point t_0 and finitely many rational exponents $\lambda_j \in \mathbb{Q}$.

The last property alone suffices to guarantee that Abelian integrals are *non-oscillating*: any such integral has only finitely many real isolated zeros on any interval. This can be proved by application of the "derivation-division algorithm" [16]. Moreover, the Fewnomial theory [3] which can be regarded as a multidimensional generalization of the above algorithm, implies the following uniform version of the non-oscillation property, achieved in [4, 18].

Theorem 1 (A. Khovanskii–A. Varchenko, 1984). The total number of real isolated zeros of Abelian integrals over the level curves of a polynomial h is uniformly bounded over all polynomial 1-forms ω and all Hamiltonians h sufficiently close to any given combination (h_0, ω_0) .

By the standard compactness and homogeneity arguments, Theorem 1 implies that for any combination of $n, m \in \mathbb{N}$ the number of real isolated zeros of Abelian integrals of forms of degree $\leq m$ over Hamiltonians of degree $\leq n$ is bounded by a constant $N_0(n, m)$ depending only on n, m. However, this constant is absolutely existential (non-constructive). Other, completely different tools allow to place *explicit* upper bounds on the number of zeros of Abelian integrals which are at least δ -distant from the critical values of the Hamiltonian [10, 11] (the bounds depend on $\delta > 0$).

One can expect similar behavior of the Poincaré–Pontryagin integrals in the case when θ is exact albeit non-polynomial anymore, $\theta = df$, with $f(x, y) = \frac{P(x,y)}{Q(x,y)}$ a rational function, though the subject was not explored systematically.

1.2.2. Pseudoabelian integrals. Completely different is the picture if the (unperturbed) closed form θ is non-exact. In this article we consider only the simplest case when θ has at most first order poles after extension on $\mathbb{C}P^2$ and hence a Darbouxian integral of the form

$$f(x,y) = \prod_{j} p_j(x,y)^{\lambda_j}, \qquad \lambda_j \in \mathbb{R}_+, \ j = 1, \dots, m,$$
(1.3)

ramified over the separatrix polygon $S(\theta) = \bigcup_j \{p_j = 0\} \subset \mathbb{R}^2$ corresponding to the "critical" value f = 0. For lack of better name, we will refer to integrals of the form

$$I(t) = \oint_{L_t} \omega, \qquad L_t \subseteq \{f = t\} \subseteq \mathbb{R}^2, \tag{1.4}$$

where ω is a rational form without singularities on L_t , as *pseudoabelian* integrals.

This particular case already contains all difficulties which distinguish the general integrable case from the Hamiltonian one. The closed level curves $\{f = t\}$ cease to be algebraic unless all ratios λ_j/λ_k are rational, moreover, their complexifications are generically dense in \mathbb{C}^2 [6]. The topological methods of continuation of the corresponding integral to the complex domain, which work so nicely in the Hamiltonian case, fail [13], thus the analytical continuation of the corresponding integral I(t) should be achieved by completely different methods. Moreover, most known approaches to the Infinitesimal Hilbert 16th problem in the Hamiltonian case use heavily the fact that I(t) satisfies some linear ordinary differential equation of finite order. This crucial property fails in the case of pseudoabelian integrals: looking on the local representation (1.5) one can see that I(t) has infinitely many linearly independent branches, which already implies that I(t) is not a solution of a Fuchsian ordinary differential equation of finite order.

Even the most basic among the properties of the Abelian integrals, the local analytic representation, fails for pseudoabelian integrals. More precisely, it can be shown that if the separatrix curve $S(\theta)$ has only normal crossings and the collection of Darboux exponents $[\lambda_1 : \cdots : \lambda_n]$ satisfies certain generic *arithmetic* properties, then the pseudoabelian integral I(t) can be locally near t = 0 represented as a composition,

$$I(t) = F(t, t^{\mu_1}, \cdots, t^{\mu_n}), \qquad \mu_1, \dots, \mu_n \in \mathbb{R}_+,$$
(1.5)

with a real analytic function F of n + 1 variables. Indeed, Lemma 3 implies this with some formal power series \hat{F} instead of F, and this formal power series converges for generic tuple of $[\lambda_1 : \cdots : \lambda_n]$ (these conditions are written explicitly in [8]). This observation, due to C. Moura [9], suffices to prove non-oscillatory behavior of pseudoabelian integrals, but under the above arithmetical restrictions on the Darboux exponents. Such form of arithmetics-conditioned finiteness is not unknown in the general theory of o-minimal structures, see [17]. Clearly, if this arithmetic restriction is indeed necessary to rule out non-oscillation of an individual pseudoabelian integral I(t), there would be almost no hope for uniform bounds for the number of zeros of pseudoabelian integrals.

1.3. **Principal result.** Our main result is the uniform non-oscillation of pseudoabelian integrals associated with a generic Darbouxian integrable foliation near the separatrix polygon $S(\theta)$. To formulate it accurately, we introduce Darbouxian classes of integrable systems. Let m_1, \ldots, m_n be a tuple of natural numbers, $m_i \in \mathbb{N}$.

Definition 1. The Darbouxian class $\mathscr{D} = \mathscr{D}(m_1, \ldots, m_n)$ is the class of real integrable foliations on $\mathbb{R}P^2$ defined by the Pfaffian equations $\theta = 0$, where θ is a *logarithmic form* which can be written in some affine chart \mathbb{R}^2 of $\mathbb{R}P^2$ in the following form³,

$$\theta = \sum_{j=1}^{n} \lambda_j \frac{dp_j}{p_j}, \qquad p_j \in \mathbb{R}[x, y], \ \deg p_j \leqslant m_j, \ \lambda_j > 0.$$
(1.6)

The polynomials p_j are assumed to be irreducible and different, hence the polar locus $S(\theta)$ is the union of n algebraic curves $\{p_j = 0\} \subset \mathbb{R}^2$.

Foliations from a fixed Darbouxian class are parameterized by points from an open subspace in the Euclidean space with the coordinates being the exponents $\lambda_1, \ldots, \lambda_n > 0$ and the respective polynomials p_j , identified with their coefficients. Using this identification, one can define open neighborhoods of a given closed form $\theta = 0$ in its Darbouxian class \mathcal{D} . We will consider only foliations which have continuous families of ovals accumulating to the separatrix $S(\theta)$ or its proper subset (as usual, in the Hausdorff sense).

If a real rational 1-form ω on \mathbb{R}^2 is *admissible* for the logarithmic form θ , i.e., if the polar locus of ω belongs to the separatrix $S(\theta)$, then the integral $I(t) = \oint_{L_t}$ over any non-singular real oval $L_t, t \neq 0$, is well defined. We claim that the number of real isolated zeros of such integrals is locally uniformly bounded.

Theorem 2 (uniform non-oscillation of pseudoabelian integrals). Let the form $\theta_0 = \sum_{j=1}^n \lambda_j \frac{dp_j}{p_j}$ defining the Darbouxian integrable foliation $\mathcal{F} = \{\theta_0 = 0\}$ be from the class $\mathcal{D}(m_1, \ldots, m_n)$, and assume that the real algebraic curves $\{p_j = 0\}$ are smooth and intersect transversally (in particular, there are no triple intersections). Let ω_0 be a rational one-form admissible for θ_0 with poles of order at most m_0 .

Then the number of isolated zeroes of the pseudoabelian integral near t = 0, as a function of (θ, ω) , is locally bounded at (θ_0, ω_0) over all admissible pairs (θ, ω) .

³One can easily show (e.g., see [6]) that a closed rational form θ with poles of order ≤ 1 on $\mathbb{C}P^2$ has the form (1.6) with complex constants λ_j .

In other words, let (θ_0, ω_0) be any admissible pair of rational 1-forms as in the Theorem. There exist a finite number $N = N(\theta_0, \omega_0)$ and $\varepsilon > 0$, depending on the pair, with the following property: for any pair (θ, ω) sufficiently close to (θ_0, ω_0) the number of real isolated zeros of any pseudoabelian integral corresponding to (θ, ω) and to a bounded family of cycles $L_t, t \in (0, \varepsilon)$, is at most N.

After this result one may expect that the number of zeros of pseudoabelian integrals would be uniformly bounded over all Darbouxian classes. The following Conjecture is a relaxed form of the Infinitesimal Hilbert 16th Problem (Conjecture 1)

Conjecture 2. For any Darbouxian class $\mathscr{D}(m_1, \ldots, m_n)$ and any degree m_0 there exists a finite number $N = N(m_0, m_1, \ldots, m_n)$ such that any pseudoabelian integral of a rational 1-form ω of degree $\leq m_0$ along ovals of the corresponding Darbouxian integral (1.3) has no more than N real isolated zeros.

Again in contrast with the Hamiltonian case, the Conjecture does not follow from Theorem 2, since the compactness arguments fail for pseudoabelian integrals. Indeed, the assumption on the separatrix is an open condition in \mathscr{D} which in the limit may degenerate into non-transversal singularities at the corners. Besides, the Darbouxian classes themselves are non-compact and in the limit contain closed forms with poles of higher orders and "vanishing separatrices" as $\lambda_j \to 0^+$. All these scenarios require additional study.

1.4. **Proofs: Mellin transform and generalized Petrov operators.** The most fundamental tool used to prove nonoscillation-type results, is the classical derivation-division algorithm [16] which goes as far back as to Descartes. In the modern language it consists in constructing a differential operator \mathcal{D} (in general, with variable coefficients) such that the result of the derivation $\mathcal{D}I(t)$ is a function without real isolated zeros (e.g., identical zero). By the Rolle-Descartes theorem, the number of zeros of I is then bounded by a constant explicitly expressible in terms of \mathcal{D} (in the simplest case of operators with constant coefficients, by the order of \mathcal{D} , provided that all characteristic roots are real). This observation was generalized in several directions, among them for the vector-valued [3] and complex analytic [12] functions.

However, this method encounters serious difficulties when applied to parametric families of functions, especially the families which admit *asymptotic* (rather than convergent) representation as in (1.5) as $t \to 0^+$, with the exponents μ_j depending on parameters and eventually coinciding between themselves. The operator \mathcal{D} , mentioned above, turns out to be very sensitive to these parameters.

Our proof is based on introducing a new class of real *pseudo*differential operators \mathcal{P} . One can define \mathcal{P} as a finite difference operator applied to the analytic continuation of a function. Remarkable fact, going back to at

least Petrov, see [14], is that application of \mathcal{P} decreases in a controlled way the number of isolated zeros. More precisely, for each such operator \mathcal{P} and any family F of functions analytic on the universal cover of the punctured closed unit disk and analytically depending on parameter there exists a finite number ("Rolle index") $N = N(\mathcal{P}, F)$ such that for any $u = u(t) \in F$ on (0, 1) the numbers of real isolated zeros of u and $\mathcal{P}u$ are related by the "Rolle inequality"

$$\#\{u=0\} \leqslant \#\{\mathcal{P}u=0\} + N.$$
(1.7)

These operators, depending on auxiliary parameters, in a sense interpolate between the "usual" differential operators, the Petrov difference operators that were used implicitly in [14] and explicitly in [15], and the operator of "taking imaginary part" which played the key role in [11].

More interestingly, the operator \mathcal{P} can be defined via the *Mellin transform* which associates with each function u(t) (say, defined on the interval [0, 1]), the function v(s) of a complex argument s by the formula

$$v = \mathcal{M}u, \qquad v(s) = \int_0^1 t^{s-1}u(t) \, dt.$$
 (1.8)

We show (Theorem 6 below) that the Mellin transform $\mathcal{M}I$ of a pseudoabelian integral I(t) is a function holomorphic in the right half-plane $\operatorname{Re} s > s_0$, which admits analytic continuation as a meromorphic function in the entire plane $s \in \mathbb{C}$ with poles located on finitely many real arithmetic progressions on the shifted negative semiaxis \mathbb{R}_- . The distance between subsequent poles may be very small, depending on the arithmetic of the exponents λ_j , which means that the residues of $\mathcal{M}I$ can be arbitrary big. Therefore attempts to rearrange the converging representation to a sum of simple fractions can result in the divergence of this rearrangement, which is another view on the divergence of the representation (1.5).

The operator P is constructed as

$$P = \mathscr{M}^{-1} \mathcal{K} \mathscr{M}, \qquad \mathcal{K} : v(s) \mapsto K(s) v(s), \tag{1.9}$$

where K(s) is an entire function $\prod_i \sin \lambda_i s$, which has zeros on these progressions. Then the product $K \cdot \mathcal{M}I$ becomes an entire function, and the only entire Mellin transform of a function from our class is identically zero. Therefore the inverse Mellin transform \mathcal{M}^{-1} is identically zero, so has no isolated zeros. Note that a differential operator with constant coefficients (the class sufficient for many applications of the derivation-division algorithm) can be represented in this way with a *polynomial* kernel K(s): this explains the connections with the traditional method.

Flexibility of the construction of the operator \mathcal{P} (location of zeros of K(s)) allows to choose $\mathcal{P} = \mathcal{P}_{\theta}$ in a way *analytically depending* on the form $\theta \in \mathscr{D}$ such that $\mathcal{P}_{\theta}I \equiv 0$ for any pseudoabelian integral $I = I_{\theta}(t)$ associated with the integrable foliation $\theta = 0$.

The cornerstone of the proof is the demonstration of the Rolle inequality (1.7) for the operator \mathcal{P} . We prove that the corresponding constant N =

 $N(\mathcal{P}_{\theta})$ depends in a uniformly bounded way on the logarithmic 1-form $\theta \in$ \mathscr{D} . This is sufficient for the proof of Theorem 2.

1.5. Note added in the proof. Since the submission of this manuscript, another, more geometric proof of Theorem 2 appeared, see [8]. In this paper authors prove that the operator \mathcal{P} annihilates the pseudoabelian integral I(t)by constructing an analytic continuation of the contour of integration into a small punctured neighborhood of zero, thus avoiding the Mellin transform entirely.

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2. Analytic representation of pseudoabelian integrals

2.1. Parametrization of admissible pairs of forms. We define the complexification $\mathscr{D}_{\mathbb{C}} = \mathscr{D}_{\mathbb{C}}(m_1, ..., m_n)$ of \mathscr{D} as the space of logarithmic forms θ ,

$$\theta = \sum_{j=1}^{n} \lambda_j \frac{dp_j}{p_j}, \qquad \lambda_j \in \mathbb{C}, \ p_j \in \mathbb{C}[x, y], \ \deg p_j \leqslant m_j,$$

where, as above, p_j are assumed to be coprime and irreducible.

Denote by \mathscr{D}' the subset $\mathscr{D}' \subset \mathscr{D} \subset \mathscr{D}_{\mathbb{C}}$ of all $\theta \in \mathscr{D}$ with smooth real algebraic curves $\{p_j = 0\}$ intersecting transversally.

The pairs (θ, ω) with $\theta \in \mathscr{D}'$ and real one-form ω admissible for θ with poles of order less than m_0 form a set $\mathscr{B} = \mathscr{B}(m_0, m_1, ..., m_n)$, which has an evident topology of a real analytic vector bundle with base \mathscr{D}' .

The set \mathscr{B} is a real analytic subvariety of the set $\mathscr{B}_{\mathbb{C}}$, defined as the set of pairs of complex one-forms (θ, ω) , with $\theta \in \mathscr{D}_{\mathbb{C}}$ and ω admissible for θ with poles of order less than m_0 . The set $\mathscr{B}_{\mathbb{C}}$ has an evident structure of a complex analytic vector bundle over complex analytic manifold $\mathscr{D}_{\mathbb{C}}$.

2.2. Local linearization of logarithmic foliations. We start with a wellknown fact whose demonstration is only briefly sketched for completeness of exposition.

Lemma 1. A closed meromorphic 1-form θ with first order poles on two transversally intersecting analytic curves in \mathbb{C}^2 can be analytically linearized near the point of their intersection. More exactly, there is a linearizing transformation ϕ mapping a neighborhood of a closed polydisk $\{|x|, |y| \leq$ $\{2\} \subset \mathbb{C}^2$ to a neighborhood of the point of the intersection, such that the pullback $\phi^*\theta$ is the logarithmic one-form $\lambda \frac{dx}{x} + \mu \frac{dy}{y}$ defined in this neighborhood. This linearizing transformation depends analytically on the complex form

 $\theta \in \mathscr{D}_{\mathbb{C}}.$

Proof. The two curves can be simultaneously rectified to become the coordinate axes. The local fundamental group of the complement is commutative, and is generated by two loops around these axes: this means that there exist two (complex) numbers λ, μ (the residues of θ) such that the difference $\theta - (\lambda \frac{dx}{x} + \mu \frac{dy}{y})$ is exact, the differential of a meromorphic function. Because of the constraints on the order of the poles, this difference is the differential of a holomorphic function dg, holomorphically depending on θ . The inverse of the holomorphic transformation $(x, y) \mapsto (x, y \exp(\mu^{-1}g))$ brings the form θ to the required "linear" normal form in some neighborhood of the origin, which can be subsequently expanded to the polydisk by a linear dilatation in x, y.

Here is the linearization transformation written explicitly in terms of the first integral. Let

$$f = Cp_1(x, y)^{\lambda} p_2(x, y)^{\mu} \prod_{j=3}^n p_j(x, y)^{\lambda_j}$$
(2.1)

be a first integral of θ , $C \neq 0$, and assume that the transversally intersecting curves are $\{p_1 = 0\}$ and $\{p_2 = 0\}$. The inverse of the mapping

$$(x,y) \to (\tilde{x},\tilde{y}), \quad \text{where } \tilde{x} = p_1(x,y), \ \tilde{y} = p_2 \prod_{j=3}^n p_j(x,y)^{\lambda_j/\mu}$$

is the required diffeomorphism ϕ .

We rescale f in order to get nice local representations near the saddles. Namely, we want f near the saddles to be equivalent, up too holomorphic change of variables to $\tilde{x}^{\lambda}\tilde{y}^{\mu}$ in the bidisc $\{|x|, |y| \leq 1\}$.

Lemma 2. For any C bigger than some $C_0 \gg 1$ the dilatation can be chosen in such a way that the first integral (2.1) becomes a monomial $\tilde{x}^{\lambda} \tilde{y}^{\mu}$ in the new coordinates.

Indeed, after linearization as in Lemma 1 the first integral becomes const $\tilde{x}^{\lambda}\tilde{y}^{\mu}$, and if const > 1, then an additional dilatation brings the first integral to the required form.

2.3. Representation for pseudoabelian integral.

2.3.1. Local computation near saddles. We compute I(t) by evaluating separately integrals along pieces δ_{ij} of L_t lying near saddles $\{p_i = p_j = 0\}$ and along the pieces δ_j of L_t lying near smooth pieces of $\{p_j = 0\}$.

Integrals of monomial 1-forms along the piece δ_{ij} of the integral curve $L_t = \{f = t\}$ lying between two cross-sections, $\{x = 1\}$ and $\{y = 1\}$, can be explicitly computed. Indeed, in the linearized coordinates (x, y) the curve δ_{ij} is, by Lemma 2, given by the explicit formula

$$\delta_{ij} = \{ y = t^{1/\mu} x^{-\lambda/\mu} \}, \qquad t \in (0,1), x \in (t^{1/\lambda}, 1), \tag{2.2}$$

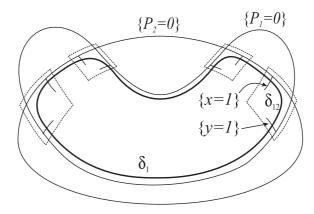


FIGURE 1. Splitting of a cycle into pieces.

eventually after replacing f by Cf, with $C \gg 1$. Therefore the integration of $\omega = x^{p-1}y^q dx$, $(p,q) \in \mathbb{Z}^2$, is simple:

$$\int_{t^{1/\lambda}}^{1} x^{p-1} y^{q} \, dx = t^{q/\mu} \left. \frac{x^{p-q\lambda/\mu}}{p-q\lambda/\mu} \right|_{t^{1/\lambda}}^{1} = \lambda^{-1} \frac{t^{q/\mu} - t^{p/\lambda}}{p/\lambda - q/\mu}, \tag{2.3}$$

if $q/\mu \neq p/\lambda$. For the resonant case $q/\mu = p/\lambda$ the same computation yields the answer $-\lambda^{-1}t^{q/\mu}\log t$ which can be otherwise obtained by direct passage to limit. Computation of $\int_{\delta_{ij}} x^p y^{q-1} dy$ is similar.

Denote by $\ell_{pq\lambda\mu}(t)$ the following elementary function:

$$\ell_{pq\lambda\mu}(t) = \begin{cases} \frac{t^{p/\lambda} - t^{q/\mu}}{p/\lambda - q/\mu}, & \text{if } p/\lambda \neq q/\mu, \\ t^{p/\lambda} \log t & \text{otherwise,} \end{cases}$$
(2.4)

One can verify that $\ell_{pq\lambda\mu}(t)$ depends analytically on λ, μ for any t > 0.

2.3.2. Analytic representation of pseudoabelian integral.

Lemma 3. After some rescaling of t, the pseudoabelian integral I(t) corresponding to (θ, ω) sufficiently close to (θ_0, ω_0) in \mathscr{B} admits the representation near the point t = 0

$$I(t) = \sum_{p,q,\lambda,\mu} a_{pq\lambda\mu} \ell_{pq\lambda\mu}(t) + \sum_{r,\lambda} b_{r\lambda} t^{r/\lambda}, \qquad (2.5)$$

where p, q, r run over a translated octant \mathbb{Z}^3_+ , $p, q, r \in -m_0 + \mathbb{N}$, and $\lambda, \mu \in \Lambda = \{\lambda_1, \ldots, \lambda_n\} \subset \mathbb{R}^n_+$ – the set of the Darboux exponents of the closed 1-form θ .

Moreover, the coefficients $a_{pq\lambda\mu}, b_{r\lambda}$ depend analytically on the forms θ, ω , and can be continued holomorphically to some neighborhood U of $(\theta_0, \omega_0) \in \mathscr{B}_{\mathbb{C}}$, and continuously to its closure $\overline{U} \subseteq \mathscr{B}_{\mathbb{C}}$. The coefficients $a_{\bullet\lambda\mu}, b_{\bullet\lambda}$ decrease at least exponentially: there exist $C, \rho > 2$ such that the following bounds hold in \overline{U}

$$|a_{pq\lambda\mu}| \leqslant C\rho^{-(p+q)}, |b_{r\lambda}| \leqslant C\rho^{-r}.$$
(2.6)

Proof. By assumptions, the Hausdorff limit $\overline{L}_0 = \lim_{t\to 0^-} L_t \subset S(\theta)$ of the real ovals is a bounded curvilinear separatrix polygon, whose edges and vertices depend analytically on $\theta \in \mathscr{D}$ in a sufficiently small neighborhood of θ_0 by the Implicit Function Theorem.

Near each vertex one can linearize the foliation using Lemma 1 in such a way that the first integral is a monomial in the new coordinates, $f = x^{\lambda}y^{\mu}$, eventually after a rescaling of f.

The linearizing transformation brings the rational form ω into a meromorphic form in the bidisk $\{|x|, |y| \leq 2\}$ with poles of order less than m_0 on the coordinate axes. Such form can be represented as a series

$$\sum_{p,q>-m_0} \alpha_{pq} x^{p-1} y^q \, dx + \beta_{pq} x^p y^{q-1} \, dy$$

converging in the bidisk $\{|x|, |y| \leq 2\}$. This convergence implies the exponentially decreasing upper bounds (2.6) on the Taylor coefficients α_{pq}, β_{pq} of the form ω in the linearizing coordinates x, y. Since the linearization depends analytically on θ in a closure $\overline{U} \in \mathscr{B}_{\mathbb{C}}$ of a sufficiently small neighborhood U of $\theta_0 \in \mathscr{B}_{\mathbb{C}}$ by Lemma 1, the coefficients α_{pq}, β_{pq} are analytic in \overline{U} and the upper bounds are uniform in this neighborhood. Integrating the series termwise along $\delta_{ij}(t) = \{f = t\} \cap \{|x|, |y| \leq 2\}$ as in (2.3), and summing up over all vertices of \overline{L}_0 corresponding to the pair (λ, μ) (their number is finite by Bezout theorem) we obtain the first sum in (2.5) and the exponential upper bounds (2.6) for $a_{pq\lambda\mu}$, as finite linear combinations of α_{pq}, β_{pq} .

The local cross-sections near the corners given in linearized coordinates by $\{x = 1\}$ and $\{y = 1\}$ become analytic curves transversal to $S(\theta)$ in the initial coordinates on the plane. The integral of the form ω along arcs of leaves L_t between the two cross-sections near two endpoints of the *j*th edge (on the curve $\{p_j = 0\}$), depends analytically on the point on the arc, as long as the arc remains nonsingular. Now, if we denote by *t* the restriction of the Darbouxian integral *f* to the cross-section, then t^{1/λ_j} will be a holomorphic parametrization of the cross-section, where λ_j is the residue (Darboux exponent) corresponding to the edge. This means that the contribution of the arc lying near this edge to the integral is an analytic function of t^{1/λ_j} . Again, after a rescaling one can assume that this analytic function converges in a small neighborhood of $\{|t^{1/\lambda_j}| \leq 2\}$, uniformly over a small neighborhood of $(\theta_0, \omega_0) \in \mathscr{B}_{\mathbb{C}}$, thus giving the second inequality in (2.6).

3. The class \Im and its Mellin transform

We introduce a class \mathfrak{J} of functions of one variable as sums of series of type (2.5) with coefficients satisfying (2.6), and define the notion of analytic

 \mathfrak{J} -family in Section 4 using the statement of Lemma 3 as its definition. We prove in Section 5 that the number of zeros on [0,1] is uniformly bounded for any analytic \mathfrak{J} -family.

3.1. Definition of the class \mathfrak{J} .

Definition 2. We define \mathfrak{J} -series as a formal series σ of the form

$$\sum_{p,q,\lambda,\mu} a_{pq\lambda\mu} \ell_{pq\lambda\mu}(t) + \sum_{r,\lambda} b_{r\lambda} t^{r/\lambda}, \qquad a_{pq\lambda\mu}, b_{r\lambda} \in \mathbb{C},$$
(3.1)

where $\lambda, \mu \in \Lambda = \{\lambda_1, ..., \lambda_n\} \subset \mathbb{R}_+$ - a finite set, $p, q, r \in -m_\sigma + \mathbb{N}$, and

$$|a_{pq\lambda\mu}| < C\rho^{-p-q}, |b_{r\lambda}| < C\rho^{-r} \quad \text{for some } C, \rho > 2.$$
(3.2)

The set Λ is called the spectrum of σ .

We denote further by $M = M(\sigma)$ some number smaller than all exponents in (3.1), e.g.:

$$M = \min(-m_{\sigma}/\lambda_i). \tag{3.3}$$

The class \mathfrak{J} is the class of functions f(t) in one variable $t \in [0, 1]$, called \mathfrak{J} -functions, representable as a sum of some \mathfrak{J} -series.

One should emphasize that the \mathfrak{J} -series corresponding to a function $f \in \mathfrak{J}$ is not unique: indeed, even the function $\ell_{pq\lambda\mu}(t)$ can be written as a sum of two monomials. Also, we do not know any effective way to check if a given function belongs to \mathfrak{J} . However, according to Lemma 3, the pseudoabelian integrals belong to \mathfrak{J} .

3.2. Analytic continuation of \mathfrak{J} -functions. We prove here that \mathfrak{J} -functions can be analytically continued to the universal cover of the punctured unit disk.

We start with an elementary inequality.

Lemma 4. For any $t, \alpha, \beta \in \mathbb{C}$ we have

$$\left|\frac{t^{\alpha} - t^{\beta}}{\alpha - \beta}\right| \leqslant |t^{\gamma}||\log t|$$

for some $\gamma \in [\alpha, \beta] \subset \mathbb{C}$.

Indeed, take $z = \log t$ in the following inequality

$$\left|e^{\alpha z} - e^{\beta z}\right| = \left|z \int_{\beta}^{\alpha} e^{wz} dw\right| \leq |\alpha - \beta| \cdot |e^{\gamma z}| \cdot |z| \quad \text{for some } \gamma \in [\beta, \alpha]. \quad \Box$$

Let $\widetilde{\mathbb{C}^*}$ denote the universal cover of $\mathbb{C}^* = \mathbb{C} \setminus \{0\}$ and let $\tilde{D}^\circ \subset \widetilde{\mathbb{C}^*}$ be the universal cover of the punctured closed unit disk.

Lemma 5. For any A > 0 the series (3.1) converges uniformly after multiplication by t^{-M} in a neighborhood in $\widetilde{\mathbb{C}^*}$ of the sector $S_A = \{|t| \leq 1, |\arg t| \leq A\} \subset \widetilde{D}^\circ$, and, in particular, on [0, 1].

Lemma 4 implies that for some constant C_A and for all p, q, λ, μ in (3.1) we have $|t^{-M}\ell_{pq\lambda\mu}(t)| \leq C_A 2^{(p+q)}$ as long as $|t|^{1/\lambda}, |t|^{1/\mu} < 2$. These inequalities define a neighborhood of S_A in which the series (3.1) converges uniformly by (3.2).

Corollary 1. \mathfrak{J} -functions are real analytic functions on (0, 1] admitting analytic continuation to the universal cover of the punctured closed unit disc.

Remark 1. If the ratios of all consecutive Darboux exponents λ_i/λ_j , $i, j = 1, \ldots, n$ are either rational, or "nice" irrational (not admitting too close approximations by rational numbers so that the differences $p\lambda_i - q\lambda_j$ decrease no faster than exponentially with $p, q \to +\infty$), then the series (2.5) can be re-expanded in powers of t^{1/λ_j} , and this re-expansion *converges*. Then the integral I(t) becomes representable under the form (1.5). Application of the standard fewnomial technique allowed C. Moura to prove that for such collections of the Darboux exponents the pseudoabelian integral I(t) is non-oscillating, moreover, that the number of isolated zeros is locally uniformly bounded over ω , provided that the Darbouxian exponents are fixed [9]. However, variation of the exponents clearly destroys the arithmetic conditions.

Remark 2. The transcendental binomials $\ell_{pq\lambda\mu}$ which in the limit converge to logarithms, first appeared at least as early as in 1951 in the work by E. Leontovich [7] on bifurcations of separatrix loop. Since then they regularly re-appear in different disguises, most recently as the so called Ecalle–Roussarie compensator.

This ubiquity can be probably explained by the fact that the function $\ell_{\varepsilon}(t) = t(t^{\varepsilon} - 1)/\varepsilon$ is in a sense the only possible Pfaffian deformation of the "logarithm" $\ell_0(t) = -t \log t$. Indeed, the graph of the latter function is an integral curve of the linear equation $\xi_0 = 0$, where $\xi_0 = t \, dy + (y - t) \, dt$ is a Pfaffian form. This linear Pfaffian equation has a singularity of the Poincaré type (resonant node) at the origin. Any (nonlinear) deformation ξ_{ε} of the form ξ_0 is necessarily analytically linearizable by the Poincaré theorem, the linearizing chart depending analytically on the parameter ε . Among integral curves of the perturbed linear system one is the graph of ℓ_{ε} .

3.3. Mellin transform of \mathfrak{J} . The (one-sided) Mellin transform of a function $u(t) \in L^1_{loc}((0,1])$ defined on the interval (0,1] is the function $v = \mathcal{M}u$ of complex variable *s* defined by the integral transform (1.8):

$$v(s) = \mathscr{M}u(s) = \int_0^1 t^{s-1}u(t)dt.$$

If the function u grows moderately as $t \to 0^+$, $|u(t)| = o(t^{-M})$ for some finite M, then the integral in (1.8) converges uniformly on compact subsets of the half-plane {Re s > M}. In particular, Mellin transforms are defined for all \mathfrak{J} -functions.

The formal Mellin transform of a \mathfrak{J} -series is defined by its action on the basic functions

$$\mathscr{M}(t^{r/\lambda}) = \frac{1}{s + r/\lambda}, \quad \mathscr{M}(\ell_{pq\lambda\mu}) = \frac{1}{(s + p/\lambda)(s + q/\mu)}$$

and extended by linearity, so the result is a formal sum of these rational functions. We prove that this formal sum converges to a meromorphic function which is an analytic continuation of the Mellin transform of the sum of the \mathfrak{J} -series.

Lemma 6. The Mellin transform $g = \mathcal{M}f$ of a function $f \in \mathfrak{J}$ given as a sum of the series (3.1) is the sum of the converging series:

$$g(s) = \sum_{p,q,\lambda,\mu} \frac{a_{pq\lambda\mu}}{(s+p/\lambda)(s+q/\mu)} + \sum_{r,\lambda} \frac{b_{r\lambda}}{s+r/\lambda},$$
(3.4)

with the same p, q, r, λ, μ and $a_{pq\lambda\mu}, b_{r\lambda}$ as in Definition 2.

In particular, g(s) extends analytically as a meromorphic function on the entire complex plane \mathbb{C} , the poles of g(s) are of at most second order, are all real and lie in the union of n arithmetic progressions of the form $\cup_{i=1}^{n} \lambda_{i}^{-1}(m_{0} - \mathbb{N}).$

Proof. The Mellin transforms of each elementary binomial $\ell_{pq\lambda\mu}$ and of $t^{r/\lambda}$ are equal to its formal Mellin transform. Termwise integration of right-hand side of (3.1) multiplied by t^{s-1} is possible due to the uniform convergence on [0, 1] of this product for $\operatorname{Re} s > -M$ by Lemma 5. Therefore we get (3.4) in { $\operatorname{Re} s > -M$ }. This double series converges to a meromorphic function on \mathbb{C} given by the same formula.

Let γ_M be the boundary of the semistrip $\{z | \operatorname{Re} z \leq -M+1, |\operatorname{Im} z| \leq -1\}$ oriented counterclockwise, with M defined in (3.3).

For any function bounded on γ_M we define an integral transform \mathcal{C}_M as:

$$\mathcal{C}_M g(t) = \frac{1}{2\pi i} \int_{\gamma} t^{-s} g(s) ds.$$
(3.5)

Assume that a meromorphic in \mathbb{C} function g is bounded on $\{|\operatorname{Im} s| = 1\}$ and has poles on the real line. Assume that all poles of g lie to the left of some real number $-M_g$. We define $\mathcal{C}g$ as $\mathcal{C}_M g$, with some $M > M_g$.

Theorem 3. The Mellin transform $\mathscr{M}f$ of any function $f \in \mathfrak{J}$ is bounded on γ_M , where $M = M(\sigma)$ for some \mathfrak{J} -series σ representing f.

The restriction of the integral transform C to \mathcal{MJ} is the inverse operator to the Mellin transform \mathcal{M} .

For $\varkappa \in \mathbb{R}$ and $f \in \mathfrak{J}$ the transform \mathcal{C} of $e^{-i\varkappa s}\mathscr{M}f(s)$ is defined and is equal to $f(e^{i\varkappa t}t)$:

$$\mathcal{C}(e^{-i\varkappa s}\mathscr{M}f(s)) = f(e^{i\varkappa}t). \tag{3.6}$$

Abusing the language we call C the inverse Mellin transform, though its domain of definition is bigger than \mathcal{MJ} .

Proof. Evidently $|s + p/\lambda| \ge 1$ on γ , so the convergence of (3.4) is uniform and $|g| \le \sum_{p,q,\lambda,\mu} |a_{pq\lambda\mu}| + \sum_{r,\lambda} |b_{r\lambda}| < \infty$ on γ . The kernel t^{-s} decreases exponentially on γ for |t| < 1, so we can integrate $t^{-s}g(s)$ termwise on γ . Calculating the residues, we see that

$$\mathcal{C}\left(\frac{1}{s+r/\lambda}\right) = t^{r/\lambda}, \quad \mathcal{C}\left(\frac{1}{(s+p/\lambda)(s+q/\mu)}\right) = \ell_{pq\lambda\mu},$$

so the Mellin transform and the inverse Mellin transform translate (3.1) to (3.4) and vice versa, i.e. are mutually inverse if restricted to \mathfrak{J} and $\mathscr{M}\mathfrak{J}$ correspondingly.

Since $e^{-i\varkappa s}$ is bounded by $e^{|\varkappa|}$ on γ , one can compute the inverse Mellin transform (3.5) termwise, again due to the uniform convergence of the series for g(s) and the exponential decay of t^{-s} . For $\ell_{pq\lambda\mu}$ and t^{α} the (3.6) can be checked immediately.

3.4. Asymptotic series for \mathfrak{J} -functions and quasianalyticity of \mathfrak{J} . Let $f \in \mathfrak{J}$ be a sum of some \mathfrak{J} -series σ . Any \mathfrak{J} -series σ can be expanded into a formal series by powers of t:

$$\check{f} = \sum_{\alpha} t^{\alpha} \left(c_{\alpha}^{-1} + c_{\alpha}^{-2} \log t \right), \qquad \alpha \in (-m_{\sigma} + \mathbb{N})\Lambda^{-1}, \tag{3.7}$$

where Λ is the spectrum of σ and Λ^{-1} denotes the set of its reciprocals. We prove below that the result of the expansion does not depend on the choice of σ and represents the asymptotic series for f. Moreover, the class \mathfrak{J} turns out to be quasianalytic at 0: any two functions having the same asymptotic series coincide.

Lemma 7. The coefficients $c_{\alpha}^{-1}, c_{\alpha}^{-2}$ are the Laurent coefficients of the Mellin transform $g = \mathcal{M}f$ of f at the pole $s = \alpha$.

The series \check{f} is an asymptotic series for f, so does not depend on σ . If $\check{f} \equiv 0$ then $f \equiv 0$. Equivalently, the only entire function in $\mathscr{M}\mathfrak{J}$ is zero.

Proof. The first claim follows because the Laurent coefficients of g are given by exactly the same formulae as c_{α}, d_{α} :

$$c_{\alpha}^{-1} = \sum_{\substack{q,\mu\\q/\mu=\alpha}} b_{q\mu} + \sum_{\substack{p,q,\lambda,\mu\\q/\mu\neq p/\lambda=\alpha}} \frac{a_{pq\lambda\mu}}{p/\lambda - q/\mu}, \quad c_{\alpha}^{-2} = \sum_{\substack{p,q,\lambda\mu\\q/\mu=p/\lambda=\alpha}} a_{pq\lambda\mu},$$

both converging due to (3.2).

Choose some $0 < \lambda_0 \leq \min \Lambda$. Each segment $[k\lambda_0, (k+1)\lambda_0), k \in \mathbb{Z}$, has at most *n* points common with $\Lambda\mathbb{Z}$, so one can choose M_k from this segment at least $\lambda_0/2n$ -distant from $\Lambda\mathbb{Z}$ (we are interested in the case of $k \to -\infty$)).

This implies that $|z + p/\lambda| \ge \lambda_0/2n$ on $\operatorname{Re} z = M_k$, so $|g| \le C$ on the boundary γ_k of the semistrip {Re $z \le M_k$, $|\operatorname{Im} z| \le 1$ }. The main point is that *C* is independent of *k*. Denote by γ'_k the boundary of the rectangular { $M_k \le \operatorname{Re} z \le -M + 1$, $|\operatorname{Im} z| \le 1$ } and let split integration along γ in the definition of *C* into a sum of two integrals, one along γ_k and another along γ'_k , see Figure 3.4.

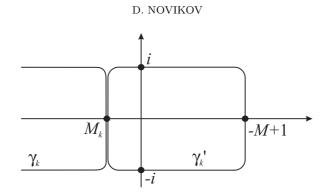


FIGURE 2. Splitting of the integration contour γ into two.

Evaluation of the integral along γ'_k gives a finite partial sum of f:

$$\frac{1}{2\pi i} \int_{\gamma'_k} t^{-s} g(s) ds = \sum_{\alpha < -M_k} t^{\alpha} \left(c_{\alpha} + d_{\alpha} \log t \right) ds$$

Integral along γ_k can be estimated from above by $2Ct^{-M_k}(1+1/\log t)$, which proves the second claim.

If $f \equiv 0$, then g has no poles at all, i.e. g is entire. Therefore the integral along γ'_k is zero for all k. Therefore $|f(t)| \leq 2Ct^{-M_k}(1+1/\log t)$ for all $t \in (0,1)$, and for all M_k . Since $M_k \to -\infty$ as $k \to -\infty$ and C is independent of k, we conclude that $f(t) \equiv 0$ on (0,1), and therefore everywhere by analyticity. This means that $g = \mathcal{M}f \equiv 0$ as well. \Box

4. Analytic $\mathscr{M}\mathfrak{J}$ -families and operator \mathcal{E}_{\varkappa}

Definition 3. Let $V \Subset \mathbb{R}^n$ be some compact set, and let $U \subset \mathbb{C}^n$ be a bounded neighborhood of V. We denote by $\mathfrak{O}_{V \subset U}$ the space of real analytic functions on V which can be extended holomorphically to U and continuously to the closure \overline{U} of U, equipped with the norm $||u||_{\mathfrak{O}} := \max_{\nu \in \overline{U}} |u(\nu|)$.

Lemma 8. The space $\mathfrak{O}_{V \subset U}$ is complete.

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Definition 4. Consider a family $\check{\mathbf{F}} = \{\sigma_{\nu}\}$ of \mathfrak{J} -series, $\nu \in V \Subset \mathbb{R}^N$, and denote by $\Lambda(\nu) = \{\lambda_1(\nu), ..., \lambda_n(\nu)\} \subset \mathbb{R}_+$ the spectrum of the \mathfrak{J} -series $\sigma(\nu)$.

We say that $\check{\mathbf{F}} = \{\sigma_{\nu}\}$ is an analytic family of \mathfrak{J} -series on pair $V \subset U$ for some bounded neighborhood $U \subset \mathbb{C}^N$ of $V \in \mathbb{R}^N$ if, first, the lower bound $m_{\sigma_{\nu}}$ defined in (3.3) can be chosen independent of ν , and, second, the functions $\lambda_i(\nu)$, $a_{pq\lambda(\nu)\mu(\nu)}(\nu)$, $b_{r\lambda(\nu)}(\nu)$ belong to $\mathfrak{D}_{V\subset U}$, the analytic continuation of $\lambda_i(\nu)$ does not vanish in \overline{U} (i.e. $\lambda_i^{-1}(\nu) \in \mathfrak{O}_{V\subset U}$ as well), and

$$\|a_{pq\lambda(\nu)\mu(\nu)}(\nu)\|_{\mathfrak{O}} \leq C\rho^{-(p+q)}, \quad \|b_{r\lambda(\nu)}(\nu)\|_{\mathfrak{O}} \leq C\rho^{-r}$$
(4.1)
for some $C = C(\check{\mathbf{F}}), \rho = \rho(\check{\mathbf{F}}) > 2.$

An analytic \mathfrak{J} -family $\mathbf{F} = \{f(t; \nu)\} \subset \mathfrak{J}$ is a family of sums of an analytic family of \mathfrak{J} -series. In other words, for every $f(t; \nu)$ one can choose a \mathfrak{J} -series σ_{ν} such that $f(t; \nu)$ is a sum of σ_{ν} , and $\{\sigma_{\nu}\}$ is an analytic family of \mathfrak{J} -series.

Analytic $\mathscr{M}\mathfrak{J}$ -family is by definition the Mellin transform of an analytic \mathfrak{J} -family.

In fact, the first requirement is obsolete, as easily follows from the following remark.

Lemma 9. Assume that the functions $\lambda_i(\nu)$, $i \in \mathbb{Z}$, are analytic in a domain $U \subset \mathbb{C}^n$, and for every $\nu \in U$ there exists $m(\nu) \in \mathbb{N}$ such that $\lambda_i(\nu) = 0$ for $i < -m(\nu)$. Then $\sup_{\nu \in U} m(\nu) < \infty$.

Indeed, the sets $S_k = \{\nu \in U | m(\nu) \leq k \in \mathbb{N}\} = \bigcap_{i < -k} \{\lambda_i(\nu) = 0\}$ are analytic subsets of U. Since $U = \bigcup_{k \in \mathbb{N}} S_k$, we have $S_k = U$ for some k. \Box

Lemma 3 says that the family of pseudoabelian integrals parameterized by \mathscr{B} is locally an analytic \mathfrak{J} -family.

Let $\dot{\mathbf{F}} = \{\sigma_{\nu}\}$ be an analytic family of \mathfrak{J} -series on pair $V \subset U$, and consider the series

$$\sum_{p,q,\lambda(\nu),\mu(\nu)} a_{pq\lambda(\nu)\mu(\nu)} \ell_{pq\lambda(\nu)\mu(\nu)}(t) + \sum_{r,\lambda(\nu)} b_{r\lambda(\nu)} t^{r/\lambda(\nu)} \quad \text{for } \nu \in U.$$
(4.2)

Lemma 10. For every A > 0 there exists some $\varepsilon > 0$ such that this series converges uniformly on compact subsets of

$$U_{\varepsilon} = \{ |\arg t| \leqslant A, 0 < |t| < e^{\varepsilon} \} \times \left(\{ |\operatorname{Im} \lambda_i^{-1}| < \varepsilon \} \cap U \right) \subset \widetilde{\mathbb{C}^*} \times U,$$

so its sum is a function holomorphic on U_{ε} .

Proof. In U_{ε} we get $|t^{\lambda^{-1}(\nu)}| \leq \exp\left(\varepsilon \left[A + \|\lambda^{-1}(\nu)\|_{\mathfrak{O}}\right]\right)$, i.e. $|\ell_{pq\lambda\mu}(t)|$ can grow exponentially fast. However, the fast decay of coefficients $a_{pq\lambda\mu}, b_{r\lambda}$ guarantees convergence if ε is sufficiently small: the series (3.1) converges uniformly on compact subsets of U_{ε} as soon as $\varepsilon \left[A + \|\lambda_i^{-1}(\nu)\|_{\mathfrak{O}}\right] < \ln 2$ for i = 1, ..., n.

Remark 3. For functions from $\mathscr{M}\mathfrak{J}$ one can try to define analytic dependence on parameter as analytic dependence of $f(t;\nu)$ on ν on compact subsets of $\mathbb{C}\setminus\mathbb{R}$. Analytic $\mathscr{M}\mathfrak{J}$ -families depend analytically on parameter in this sense, as shown above, but the opposite implication seems to be wrong.

4.1. The operator \mathcal{E}_{\varkappa} . From Lemma 11 below follows that the rotation of argument preserves \mathfrak{J} : for any $f(t) \in \mathfrak{J}$ the function $f(e^{i\varkappa}t), \varkappa \in \mathbb{R}$, is also in \mathfrak{J} . From (3.6) it seems that the Mellin conjugate to the rotation is the operator of multiplication by $e^{-i\varkappa s}$. However, the latter does not preserves $\mathscr{M}\mathfrak{J}$: the function $e^{-i\varkappa s}\mathscr{M}f(s)$ grows as $|\operatorname{Im} s| \to \infty$, while any function in $\mathscr{M}\mathfrak{J}$ decreases as $|\operatorname{Im} s|^{-1}$.

In this section we define the Mellin conjugate \mathcal{E}_{\varkappa} of the rotation on $\mathscr{M}\mathfrak{J}$, and prove that application of operator \mathcal{E}_{\varkappa} with \varkappa analytically depending on parameters preserves $\mathscr{M}\mathfrak{J}$ -families. Let us start from the standard computation applied to $f \in \mathfrak{J}$:

$$\mathcal{M}(f(e^{i\varkappa}t)) = \int_0^1 t^{s-1} f(e^{i\varkappa}t) dt$$
$$= \int_0^{e^{-i\varkappa}} t^{s-1} f(e^{i\varkappa}t) dt + \int_{C_{\varkappa}} t^{s-1} f(e^{i\varkappa}t) dt$$
$$= e^{-i\varkappa s} \int_0^1 u^{s-1} f(u) du + R(s) = e^{-i\varkappa s} \mathcal{M}f(s) + R(s), \quad (4.3)$$

where $C_{\varkappa} = \{e^{iz}, z \in [-\varkappa, 0]\}$ and $R(s) = \int_{C_{\varkappa}} t^{s-1} f(t) dt$ is an entire function of exponential type. One can immediately see that R(s) is bounded in $\{|\operatorname{Im} s| \leq 1\}$. Therefore $\mathcal{C}R \equiv 0$, which explains (3.6). We take this as a definition of \mathcal{E}_{\varkappa} . In a sense \mathcal{E}_{\varkappa} is an operator of projection of $e^{-i\varkappa s}\mathscr{M}\mathfrak{J}$ to $\mathscr{M}\mathfrak{J}$ along the space of entire functions.

Definition 5. Let $\varkappa \in \mathbb{R}$. For any $g \in \mathscr{M}\mathfrak{J}$ define $\mathscr{E}_{\varkappa}g$ as the only function in $\mathscr{M}\mathfrak{J}$ such that the difference $e^{-i\varkappa s}g - \mathscr{E}_{\varkappa}(g)$ is an entire function.

The uniqueness follows from Lemma 7: if $g_1, g_2 \in \mathscr{M}\mathfrak{J}$ are two functions which differ from $e^{-i\varkappa s}g$ by entire functions then their difference is entire and lies in $\mathscr{M}\mathfrak{J}$. Then Lemma 7 implies that the difference is zero.

Remark 4. One can show that $\mathcal{E}_{\varkappa}g(s)$ is an analytic continuation from $\{\operatorname{Im} s > 1\}$ of the integral

$$\frac{1}{2\pi i} \int_{\gamma} \frac{e^{-i\varkappa z}g}{z-s} dz,$$

where γ is as in (3.5).

Lemma 11. Let $F = \{g(s; \nu)\}$ be an analytic $\mathscr{M}\mathfrak{J}$ -family on pair $V \subset U$, and let $\varkappa = \varkappa(\nu) \in \mathfrak{D}_{V \subset U}$. Then

- (1) $\{\mathcal{E}_{\varkappa(\nu)}g(s;\nu)\}\$ is an analytic $\mathscr{M}\mathfrak{J}$ -family on pair $V \subset \widetilde{U}\$ for some $\widetilde{U} \subset U.$
- (2) the operator $\mathcal{C} \circ \mathcal{E}_{\varkappa} \circ \mathscr{M}$ maps analytic \mathfrak{J} -families to analytic \mathfrak{J} -families,
- (3) $\mathcal{CE}_{\varkappa}\mathcal{M}f = f(e^{i\varkappa}t)$, and
- (4) the operator of rotation of argument $f(t) \to f(e^{i\varkappa}t)$ preserves \mathfrak{J} .

Proof. Let us describe explicitly the \mathcal{E}_{\varkappa} . This is simple for $\frac{1}{(s+p/\lambda)(s+q/\mu)}$ and $\frac{1}{s+r/\lambda}$:

$$\frac{e^{-i\varkappa s}}{s-x} = \frac{e^{-i\varkappa x}}{s-x} + \frac{e^{-i\varkappa s} - e^{-i\varkappa x}}{s-x},$$

and

$$2\frac{e^{-i\varkappa s}}{(s-x)(s-y)} = \frac{e^{-i\varkappa x} + e^{-i\varkappa y}}{(s-x)(s-y)} + \frac{e^{-i\varkappa x} - e^{-i\varkappa y}}{x-y} \left(\frac{1}{s-x} + \frac{1}{s-y}\right) + R,$$

where $R = 2 \frac{(s-x)e^{-i \varkappa y} + (x-y)e^{-i \varkappa s} + (y-s)e^{-i \varkappa x}}{(s-x)(s-y)(x-y)}$ is entire in s. We extend this by linearity to any $g \in \mathscr{M}\mathfrak{J}$. Namely, for g given by (3.4)

We extend this by linearity to any $g \in \mathcal{MJ}$. Namely, for g given by (3.4) we define $\mathcal{E}_{\varkappa}g(s)$

$$\mathcal{E}_{\varkappa}g(s) = \sum_{p,q,\lambda,\mu} \frac{\tilde{a}_{pq\lambda\mu}}{(s+p/\lambda)(s+q/\mu)} + \sum_{r,\lambda} \frac{b_{r\lambda}}{s+r/\lambda},$$

by the following formulae:

$$\tilde{a}_{pq\lambda\mu} = \frac{1}{2} \left(e^{-i\varkappa p/\lambda} + e^{-i\varkappa q/\mu} \right) a_{pq\lambda\mu}, \tag{4.4}$$

$$\tilde{b}_{r\lambda} = e^{-i\varkappa r/\lambda}b_{r\lambda} + \sum_{q,\mu} \frac{e^{-i\varkappa r/\lambda} - e^{-i\varkappa q/\mu}}{r/\lambda - q/\mu}a_{rq\lambda\mu}.$$
(4.5)

Take ε so small that $\tilde{\rho} = \rho e^{-\varepsilon} > 2$. Let a neighborhood $\widetilde{U} \subset U$ of V be so small that $\left| \operatorname{Im} \left(\varkappa(\nu) \lambda_i^{-1}(\nu) \right) \right| < \varepsilon$ for i = 1, ..., n and $\nu \in \widetilde{U}$. Then $|e^{-i\varkappa r/\lambda}| < e^{\varepsilon r}$ in \widetilde{U} .

For p, q, r > -m as in (3.4) we have $\max(p, q) , so for <math>\nu \in \widetilde{U}$ we have estimates

$$\frac{1}{2} \left| e^{-i\varkappa p/\lambda} + e^{-i\varkappa q/\mu} \right| \leqslant e^{\varepsilon(p+q+m)}, \quad \left| \frac{e^{-i\varkappa r/\lambda} - e^{-i\varkappa q/\mu}}{r/\lambda - q/\mu} \right| \leqslant e^{\varepsilon(r+q+m)} \|\varkappa\|_{\mathfrak{O}},$$

by Lemma 4.

The upper bounds (3.2) then imply that (4.4), (4.5) converge in $\mathfrak{O}_{V\subset \widetilde{U}},$ and

$$\tilde{a}_{pq\lambda\mu} \leqslant \tilde{C}\tilde{\rho}^{-p-q}, \quad \tilde{b}_{r\lambda} \leqslant \tilde{C}\tilde{\rho}^{-r} \quad \text{in } \tilde{U},$$

which proves the first claim of the Lemma. The second claim is a consequence of the first and of definition of $\mathscr{M}\mathfrak{J}$ -family. The third claim follows from (4.3) and (3.6), so the fourth follows from the third and the fact that \mathcal{E}_{\varkappa} preserves $\mathscr{M}\mathfrak{J}$.

5. Petrov operator and proof of Theorem 2

We generalize the so-called *Petrov operator*, see [15], as follows:

$$\mathcal{P}_{\varkappa}f(t) := \frac{1}{2i} \left(f(te^{-i\varkappa}) - f(te^{i\varkappa}) \right) \quad \varkappa \in \mathbb{R}$$
(5.1)

Lemma 12. Let $g = \mathcal{M}f$. The inverse Mellin transform of $g(s)\sin(\varkappa s)$ is $\mathcal{P}_{\varkappa}f$. For any $f \in \mathfrak{J}$

$$\mathcal{P}_{\varkappa}f = \mathscr{M}^{-1} \circ (\mathcal{E}_{\varkappa} - \mathcal{E}_{-\varkappa}) \circ \mathscr{M}f$$

This is a straightforward applications of (3.6) and of Lemma 11.

The key idea of the proof of Theorem 2 is a generalization of the Petrov trick invented in [14], where it was used for $\varkappa = \pi$.

Let f(t) be a \mathfrak{J} -function. Denote by N(f) the number of isolated zeroes of function f on (0,1) counted with multiplicities, by $\Delta^1_{\varkappa}(f)$ the increment of $\arg f(t)$ along the arc $\{e^{i\varkappa\phi}, \phi \in [-1,1]\}$ and let $\Delta^0_{\varkappa}(f)$ be equal to the limit as $\varepsilon \to 0$ of the increment of $\arg f(t)$ along the arc $\{\varepsilon e^{i\varkappa\phi}, \phi \in [-1, 1]\}$. Finiteness of all these numbers follows from Lemma 7.

Lemma 13. For $f \in \mathfrak{J}$ and real on (0,1) the following inequality holds:

$$N(f) \leq 1 + N(\mathcal{P}_{\varkappa}f) + \frac{1}{2\pi} (\Delta^{1}_{\varkappa}(f) + \Delta^{0}_{\varkappa}(f)).$$
(5.2)

Proof. Applying the argument principle to the sector $\{|\arg z| \leq \varkappa, \varepsilon \leq |z| \leq 1\}$ on the universal cover of the punctured closed unit disc we get

$$2\pi N(f) \leq \Delta_{-} + \Delta_{+} + \Delta_{\varkappa}^{1}(f) + \Delta_{\varkappa}^{0}(f), \qquad (5.3)$$

where Δ_{\pm} are the increments of the argument of f(t) along the rays $I_{\pm} = \{te^{\pm i\varkappa}, t \in [\varepsilon, 1]\}$. Since f(t) is real on the real axis, it takes conjugate values on I_{\pm} , so $\Delta_{+} = \Delta_{-}$ and $\operatorname{Im} f|_{I_{\pm}} = \pm \mathcal{P}_{\varkappa} f$. First, assume that $\mathcal{P}_{\varkappa} f \neq 0$. Evidently, any segment of I_{+} where incre-

First, assume that $\mathcal{P}_{\varkappa}f \neq 0$. Evidently, any segment of I_+ where increment of argument of f(t) is greater than π contains at least one zero of Im f(t), so

$$\Delta_{+} \leqslant \pi \left(N(\operatorname{Im} f|_{I_{+}} + 1) = \pi N(\mathcal{P}_{\varkappa}f) + \pi_{*} \right)$$

and we get the (5.2).

If $\mathcal{P}_{\varkappa}f \equiv 0$, then, by symmetry, $f(te^{i\varkappa}) = f(te^{-i\varkappa})$, i.e. $f = \tilde{f}(t^{\pi/\varkappa})$ with \tilde{f} holomorphic in a punctured unit disc. Therefore (5.2) follows from the argument principle applied to \tilde{f} and a punctured unit disc: in this case we do not need I_{\pm} , and the contour consists of two circles of radius 1 and ε . \Box

Remark 5. As $\varkappa \to 0$, $\mathcal{P}_{\varkappa}f(t)$ tends to a multiple of f'(t), the contributions $\Delta^1_{\varkappa}(f)$ and $\Delta^0_{\varkappa}(f)$ of arcs tend to zero, and we get the standard Rolle theorem as a limit case of Lemma 13.

Let $\mathbf{F} = \{f(t, \nu)\}$ be a \mathfrak{J} -family on pair $V \subset U$, and let $\varkappa = \varkappa(\nu) \in \mathfrak{O}_{V \subset U}$.

Lemma 14. The variation $\Delta^1_{\varkappa(\nu)} f$ is uniformly bounded from above by some $\Delta^1 = \Delta^1(\mathbf{F}, \varkappa) < \infty$ for all $\nu \in V$.

Proof. The functions $\operatorname{Im} f(e^{i \varkappa(\nu)\phi}; \nu)$ and $\operatorname{Re} f(e^{i \varkappa(\nu)\phi}; \nu)$ are analytic functions of $\phi \in [-1, 1]$ analytically depending on parameter $\nu \in V \Subset \mathbb{R}^n$. Therefore their number of isolated zeros counted with multiplicity is uniformly bounded from above by some number by Gabrielov's theorem [2]. The variation Δ^1 is at most π (this number + 1). \Box

Lemma 15. The variation $\Delta^0_{\varkappa(\nu)} f$ is uniformly bounded from above by some $\Delta^0 = \Delta^0(\mathbf{F}, \varkappa) < \infty$ for all $\nu \in V$.

Proof. For a sum $f \in \mathfrak{J}$ of a \mathfrak{J} -series σ the number $\Delta^0_{\varkappa}(f)$ is bounded from above by $m_{\sigma}\varkappa$ by Lemma 7. Therefore the existence of the uniform upper bound follows from the Lemma 9.

Corollary 2. For F, \varkappa as above and for any $f \in F$ $N(f) \leq N(\mathcal{P}_{\varkappa}f) + \Delta(F, \varkappa), \quad \Delta(F, \varkappa) = \Delta^{0}(F, \varkappa) + \Delta^{1}(F, \varkappa) + 1 < \infty.$ (5.4)

5.1. **Proof of Theorem 2.** Starting from the analytic \mathfrak{J} -family F of pseudoabelian integrals parameterized by $(\theta, \omega) \in U \subset \mathscr{B}$ we construct the new family $\{\mathcal{P}_{\pi\lambda_n(\nu)}f(\bullet,\nu)\}$. It is again an analytic \mathfrak{J} -family by Lemma 11. By Corollary 2, application of the operator $\mathcal{P}_{\pi\lambda_n(\nu)}$ reduces the number of isolated zeros at most by a finite number uniformly over parameters.

On the other hand, the new family is simpler: its Mellin transform now has poles on $\{m/\lambda_k(\nu), k = 1, ..., n-1, m \in \mathbb{Z}_{>-m_f}\}$, i.e. on the union of only n-1 arithmetic progressions. Indeed, the Mellin transform of $\{\mathcal{P}_{\pi\lambda_n(\nu)}f_\nu\}$ has the same poles as $\sin(\pi\lambda_n(\nu)s)\mathcal{M}f$, and the first factor has simple zeros at $\{m/\lambda_n(\nu), m \in \mathbb{Z}_{>-m_f}\}$. This means that the simple poles of $\mathcal{M}f$ on this progression will become regular points of $\mathcal{MP}_{\pi\lambda_n(\nu)}f$, and double poles, which belong to two arithmetic progression simultaneously, will become simple poles.

Repeating this step n times we arrive at the analytic \mathfrak{J} -family whose Mellin transform consists of entire functions. This means that this analytic \mathfrak{J} -family is identical zero by Lemma 7, so has no isolated zeros. \Box .

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